

Construction of quasi-periodic response solutions in forced strongly dissipative systems

Guido Gentile

Dipartimento di Matematica, Università di Roma Tre, Roma, I-00146, Italy.

E-mail: gentile@mat.uniroma3.it

Abstract

We consider a class of ordinary differential equations describing one-dimensional quasi-periodically forced systems in the presence of large damping. We give a fully constructive proof of the existence of response solutions, that is quasi-periodic solutions which have the same frequency vector as the forcing. This requires dealing with a degenerate implicit function equation: we prove that the latter has a unique solution, which can be explicitly determined. As a by-product we obtain an explicit estimate of the minimal size of the damping coefficient.

1 Introduction

In this paper we continue the analysis started in [9, 10, 8] on the existence and properties of quasi-periodic motions in one-dimensional strongly dissipative forced systems.

We consider one-dimensional systems with a quasi-periodic forcing term in the presence of strong damping, described by ordinary differential equations of the form

$$\varepsilon \ddot{x} + \dot{x} + \varepsilon g(x) = \varepsilon f(\omega t), \quad (1.1)$$

where $g: \mathbb{R} \rightarrow \mathbb{R}$ and $f: \mathbb{T}^d \rightarrow \mathbb{R}$ are real analytic functions and $\mathbb{T} = \mathbb{R}/2\pi\mathbb{Z}$. We call $g(x)$ the *mechanical force*, $f(\omega t)$ the *forcing term*, $\omega \in \mathbb{R}^d$ the *frequency vector* of the forcing, and $\gamma = 1/\varepsilon > 0$ the *damping coefficient*.

The function f is quasi-periodic in t , i.e.

$$f(\psi) = \sum_{\nu \in \mathbb{Z}^d} e^{i\nu \cdot \psi} f_\nu, \quad \psi \in \mathbb{T}^d, \quad (1.2)$$

with average $\langle f \rangle = f_0$, and \cdot denoting the scalar product in \mathbb{R}^d . By the analyticity assumption on f and g , one has $|f_\nu| \leq \Phi e^{-\xi|\nu|}$ and $|d^s g(c_0)/dx^s| \leq s! \Gamma^s$ for suitable positive constants Φ , ξ , and Γ .

A Diophantine condition is assumed on ω . Define the *Bryuno function*

$$\mathfrak{B}(\omega) = \sum_{n=0}^{\infty} \frac{1}{2^n} \log \frac{1}{\alpha_n(\omega)}, \quad \alpha_n(\omega) = \inf\{|\omega \cdot \nu| : \nu \in \mathbb{Z}^d \text{ such that } 0 < |\nu| \leq 2^n\}. \quad (1.3)$$

Assumption 1. *The frequency vector ω satisfies the Bryuno condition $\mathfrak{B}(\omega) < \infty$.*

The following assumption will be made on the functions g and f (for given force $g(x)$ this can read as a condition on the forcing term).

Assumption 2. *There exists $c_0 \in \mathbb{R}$ such that $x = c_0$ is a zero of odd order \mathbf{n} of the equation*

$$g(x) - f_0 = 0, \tag{1.4}$$

that is $g_0 = \mathbf{n}!^{-1}d^n g/dx^n(c_0) \neq 0$ and, if $\mathbf{n} > 1$, $d^k g/dx^k(c_0) = 0$ for $k = 1, \dots, \mathbf{n} - 1$.

In [8] we proved the following result about the existence of quasi-periodic solutions with the same frequency vector ω as the forcing (response solutions).

Theorem 1.1. *Under Assumptions 1 and 2, for ε small enough there exists at least one quasi-periodic solution $x_0(t) = c_0 + u(\omega t, \varepsilon)$, with frequency vector ω , reducing to c_0 as ε tends to 0. Such a solution is analytic in t .*

Note that the condition $\varepsilon > 0$ could be eliminated: indeed, the proof in [8] works for all ε small enough, and the request $\varepsilon > 0$ only aims to interpret $\gamma = 1/\varepsilon$ as the damping coefficient. Analyticity on t is not explicitly stated in [8], but follows immediately from the proof therein.

We also proved in [8] that the condition that c_0 be a zero of odd order of (1.4) is a necessary and sufficient condition for a quasi-periodic solution around c_0 to exist.

However, as pointed out in [8], except for the non-degenerate case $\mathbf{n} = 1$ (where the implicit function theorem applies), in general the proof ultimately relies on continuity arguments, which do not provide a quantitative constructive estimate on the maximal size of the perturbation parameter ε . Moreover, the quasi-periodic solution was constructed in terms of two parameters, that is ε and c : the latter is defined as the constant part of the quasi-periodic solution itself and is fixed in terms of ε so as to solve a certain implicit function equation (the so-called bifurcation equation). In particular, a quasi-periodic solution was showed to exist for any solution $c(\varepsilon)$ to the bifurcation equation, but the problem of studying how many such solutions exist and how do they depend on ε was not investigated.

In this paper we show that for all odd \mathbf{n} the bifurcation equation admits one and only one solution. Furthermore, we explicitly construct the quasi-periodic solution in terms of the only parameter ε , and we also give a quantitative estimate on the maximal size of ε . So, with respect to [8], the proof of existence of the quasi-periodic solution is fully constructive.

We can summarise our results in the following statement.

Theorem 1.2. *Under Assumptions 1 to 2 on the ordinary differential equation (1.1), there exists an (explicitly computable) constant $\varepsilon_0 > 0$ such that for all $|\varepsilon| < \varepsilon_0$ there exists a quasi-periodic solution $x_0(t) = c_0 + u(\omega t, \varepsilon)$, with frequency vector ω , such that $u = O(\varepsilon)$ as $\varepsilon \rightarrow 0$. Such a solution is analytic in t and depends C^∞ -smoothly on ε .*

The result solves a problem left as open in [8]. The problem remains whether other quasi-periodic solutions exist.

Numerical algorithms to construct response solutions of quasi-periodically forced dissipative systems are provided in [3, 11, 12], based either on a generalised harmonic balance method or on a fixed point method for a suitable Poincaré map. Also the method described in the present paper is well suited for numerical implementations, and allows a completely rigorous control of the approximation error for the solution. Moreover it applies also in degenerate cases where one cannot apply directly the implicit function theorem (of course it requires for the damping coefficient to be large enough).

2 The bifurcation equation

The existence of quasi-periodic solutions $x_0(t)$ has been proved in [8]. The proof proceeds as follows. First of all, write

$$x_0(t) = c_0 + u(\omega t, \varepsilon) = c + X(\omega t; \varepsilon, c), \quad X(\psi; \varepsilon, c) = \sum_{\nu \in \mathbb{Z}_*^d} e^{i\nu \cdot \psi} X_\nu(\varepsilon, c), \quad (2.1)$$

with $\mathbb{Z}_*^d = \mathbb{Z}^d \setminus \{\mathbf{0}\}$ (so that c is the average of x_0 on the d -dimensional torus). The Fourier coefficients $X_\nu = X_\nu(\varepsilon, c)$ of the function $X(\psi; \varepsilon, c)$ are obtained by solving the *range equation*

$$i\omega \cdot \nu (1 + i\varepsilon\omega \cdot \nu) X_\nu + \varepsilon [g(c + X(\cdot; \varepsilon, c))]_\nu = \varepsilon f_\nu, \quad \nu \neq \mathbf{0}, \quad (2.2)$$

where

$$[g(c + X(\cdot; \varepsilon, c))]_\nu = \sum_{p=0}^{\infty} \frac{1}{p!} \frac{d^p}{dx^p} g(c) \sum_{\substack{\nu_1, \dots, \nu_p \in \mathbb{Z}_*^d \\ \nu_1 + \dots + \nu_p = \nu}} X_{\nu_1}(\varepsilon, c) \dots X_{\nu_p}(\varepsilon, c). \quad (2.3)$$

The analysis in [8] shows that for all c close enough to c_0 and all ε small enough there exists a function $X(\psi; \varepsilon, c)$ which solves (2.2). Moreover the map $(\varepsilon, c) \mapsto X(\cdot; \varepsilon, c)$ is C^∞ in a neighbourhood of $(0, c_0)$, and

$$X(\psi; \varepsilon, c) = \varepsilon X^{(1)}(\psi) + o(\varepsilon), \quad \dot{X}^{(1)}(\omega t) = f(\omega t). \quad (2.4)$$

Then, for any solution $c = c(\varepsilon)$ to the *bifurcation equation*

$$F(\varepsilon, c) := [g(c + X(\cdot; \varepsilon, c))]_{\mathbf{0}} - f_{\mathbf{0}} = 0, \quad (2.5)$$

there exists a quasi-periodic solution $x_0(t) = c(\varepsilon) + X(\omega t; \varepsilon, c(\varepsilon))$.

In principle, there could be several solutions to (2.5) when $n > 1$. On the contrary, we shall prove that the solution to (2.5) is unique. Unfortunately, this does not implies that the response solution $c_0 + u(\omega t, \varepsilon)$ of Theorem 1.1 is the only quasi-periodic solution reducing to c_0 as $\varepsilon \rightarrow 0$, because no result ensures that the function $X(\psi; \varepsilon, c)$ which solves (2.2) at fixed c and ε is unique.

We shall write $c = c(\varepsilon) = c_0 + \zeta$, with $\zeta = \zeta(\varepsilon)$ such that $\zeta(0) = 0$, so that (2.3) becomes

$$[g(c + X(\cdot; \varepsilon, c))]_\nu = \sum_{p=n}^{\infty} \frac{1}{p!} \frac{d^p}{dx^p} g(c_0) \sum_{k=0}^p \binom{p}{k} \zeta^{p-k} \sum_{\substack{\nu_1, \dots, \nu_k \in \mathbb{Z}_*^d \\ \nu_1 + \dots + \nu_k = \nu}} X_{\nu_1}(\varepsilon, c) \dots X_{\nu_k}(\varepsilon, c). \quad (2.6)$$

Lemma 2.1. *Let $c = c(\varepsilon)$ a solution of the bifurcation equation (2.5). Then $c = c_0 + O(\varepsilon)$, and $c = c_0 + o(\varepsilon)$ requires $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} = 0$.*

Proof. Assume that $\zeta/\varepsilon \rightarrow \infty$ as $\varepsilon \rightarrow 0$. Then, by using (2.4), one has

$$[g(c + X(\cdot; \varepsilon, c))_{\mathbf{0}} - f_{\mathbf{0}}] = g_0 [(\zeta + X(\cdot, \varepsilon, c))^n]_{\mathbf{0}} + O(\zeta^{n+1}) = g_0 \zeta^n + O(\zeta^{n+1}) = 0,$$

which leads to a contradiction. Hence $\zeta = O(\varepsilon)$. On the other hand if $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} \neq 0$ and $\zeta = o(\varepsilon)$ one has

$$[g(c + X(\cdot; \varepsilon, c))_{\mathbf{0}} - f_{\mathbf{0}}] = g_0 [(\zeta + X(\cdot, \varepsilon, c))^n]_{\mathbf{0}} + O(\varepsilon^{n+1}) = g_0 \varepsilon^n [(X^{(1)}(\cdot))^n]_{\mathbf{0}} + O(\varepsilon^{n+1}) = 0,$$

which once more leads to a contradiction. ■

Note that a sufficient condition for $[(X^{(1)}(\cdot))^n]_{\mathbf{0}}$ to vanish is that f in (1.2) is even.

The function $F(\varepsilon, c)$, defined in (2.5), is C^∞ in both ε and c ; see [8]. In order to identify the leading orders to $F(\varepsilon, c)$, we consider the carrier

$$\Delta(F) = \{(k, j) \in \mathbb{Z}_+ \times \mathbb{Z}_+ : F_{k,j} \neq 0\}, \quad F_{k,j} = \frac{d^k}{d\varepsilon^k} \frac{d^j}{dc^j} F(0, c_0), \quad (2.7)$$

and draw the Newton polygon in the (k, j) plane; see [2] (see also [4]). If we denote by $\{(k_1, j_1), (k_2, j_2)\}$ the segment with endpoints (k_1, j_1) and (k_2, j_2) in the (k, j) plane, by construction for $n > 1$ the Newton polygon consists of (cf. figure 1)

1. either the only segment $\{(0, n), (n, 0)\}$, if $F_{n,0} = [(X^{(1)}(\cdot))^n]_{\mathbf{0}} \neq 0$,
2. or the only segment $\{(0, n), (n-1, 1)\}$, if $F_{k,0} = 0$ for all $k \geq 0$,
3. or the two segments $\{(0, n), (n-1, 1)\}$ and $\{(n-1, 1), (p, 0)\}$, with $p \geq n+1$, if $F_{k,0} = 0$ for all $k \leq p-1$ and $F_{p,0} \neq 0$.

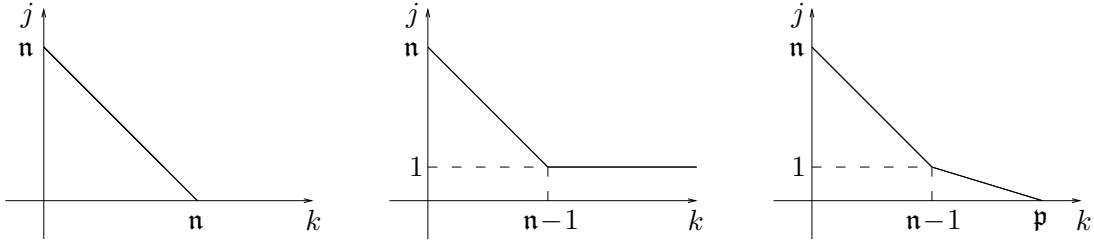


Figure 1: Newton polygon (corresponding to cases 1 to 3, respectively)

Therefore, in principle, a solution to the bifurcation equation $F(\varepsilon, c) = 0$ is

- either $c = c_0 + \zeta_1 \varepsilon + o(\varepsilon)$ for some $\zeta_1 \neq 0$, corresponding to the segment $\{(0, n), (n, 0)\}$ if $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} \neq 0$ (case 1) and to the segment $\{(0, n), (n-1, 1)\}$ if $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} = 0$ (cases 2 and 3),

- or $c = c_0 + \zeta_0 \varepsilon^{\mathbf{p}-\mathbf{n}+1} + o(\varepsilon^{\mathbf{p}-\mathbf{n}+1})$ for some $\zeta_0 \neq 0$, corresponding to the segment $\{(\mathbf{n} - 1, 1), (\mathbf{p}, 0)\}$ (case 3),
- or $c = c_0 + \zeta(\varepsilon)$, with $\zeta(\varepsilon)$ decaying to zero faster than any power of ε (case 2).

The case $\mathbf{n} = 1$ should be discussed apart, but since it is much easier and moreover has already been discussed in [9, 10], here we concentrate on $\mathbf{n} > 1$.

Lemma 2.2. *If $[(X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} = 0$ then the bifurcation equation (2.5) admits one and only one solution $c(\varepsilon) = o(\varepsilon)$. Moreover such a solution is smooth in ε , and it can be explicitly computed.*

Proof. If $[(X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} = 0$ one has

$$\begin{aligned} [(\zeta + \varepsilon X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} &= [(\zeta + \varepsilon X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} - [(\varepsilon X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} = [(\zeta + \varepsilon X^{(1)}(\cdot))^{\mathbf{n}} - (\varepsilon X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} \\ &= \mathbf{n}\zeta \left[\int_0^1 ds (\varepsilon X^{(1)}(\cdot) + s\zeta)^{\mathbf{n}-1} \right]_{\mathbf{0}} = \mathbf{n}\zeta \int_0^1 ds \left[(\varepsilon X^{(1)}(\cdot) + s\zeta)^{\mathbf{n}-1} \right]_{\mathbf{0}}, \end{aligned}$$

which is non-zero because $\mathbf{n} - 1$ is even. Therefore, by using that $\zeta = O(\varepsilon)$ by Lemma 2.1, one obtains

$$F(\varepsilon, c) = \zeta \varepsilon^{\mathbf{n}-1} a + G(\varepsilon, c), \quad a = g_0 \mathbf{n} \int_0^1 ds \left[(X^{(1)}(\cdot) + s\zeta \varepsilon^{-1})^{\mathbf{n}-1} \right]_{\mathbf{0}},$$

with $a > 0$ and $G(\varepsilon, c) = O(\varepsilon^{\mathbf{n}+1})$. Furthermore either all derivatives of $G(\varepsilon, c_0)$ vanish at $\varepsilon = 0$ (and hence $F_{k,0} = 0$ for all $k \geq 0$) or

$$G(\varepsilon, c_0) = b\varepsilon^{\mathbf{p}} + O(\varepsilon^{\mathbf{p}+1}), \quad b = F_{\mathbf{p},0},$$

for some $\mathbf{p} > \mathbf{n}$.

In both cases, there is no solution corresponding to the segment $\{(0, \mathbf{n}), (\mathbf{n} - 1, 1)\}$. Indeed, the bifurcation equation can be written as

$$\varepsilon^{\mathbf{n}-1} (a_0 \zeta + \Gamma(\varepsilon, \zeta)) = 0, \quad a_0 = \left[(X^{(1)}(\cdot))^{\mathbf{n}-1} \right]_{\mathbf{0}},$$

with the function $\Gamma(\varepsilon, \zeta)$ which is C^∞ in both ε and ζ . Hence we can apply the implicit function theorem to deduce that for $\varepsilon \neq 0$ there is one and only one solution $c = c_0 + \zeta(\varepsilon)$ to the bifurcation equation, with $\zeta(\varepsilon)$ smooth in ε and such that $\zeta(\varepsilon) = G(\varepsilon, c_0)/a_0 \varepsilon^{\mathbf{n}-1} + \tilde{\zeta}(\varepsilon)$ and $\tilde{\zeta}(\varepsilon)/\zeta(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$. Moreover one can explicitly estimate the interval U around $\varepsilon = 0$ in which such a solution exists (again by using the implicit function theorem and the bounds in [8]). \blacksquare

If $F_{k,0} = 0$ for all $k \geq 0$ then the only solution to the bifurcation equation is $c = c_0 + \zeta(\varepsilon)$, where $\zeta(\varepsilon)$ is a C^∞ function which goes to zero faster than any power as ε tends to 0. If there exists $\mathbf{p} \geq \mathbf{n} + 1$ such that $F_{k,0} = 0$ for $k \leq \mathbf{p}$ and $F_{\mathbf{p},0} \neq 0$, the only solution is of the form $c = c_0 + \zeta_0 \varepsilon^{\mathbf{p}-\mathbf{n}+1} + o(\varepsilon^{\mathbf{p}-\mathbf{n}+1})$, where ζ_0 is the (unique) solution of the equation $a\zeta + b = 0$.

In the case $[(X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} \neq 0$ the following result holds.

Lemma 2.3. *If $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} \neq 0$ then all solutions of the bifurcation equation (2.5) are of order ε and have multiplicity 1. In particular, there exists a constant $a_0 > 0$ such that $|c_i(\varepsilon) - c_j(\varepsilon)| > a_0|\varepsilon|$ for all pairs of such solutions $c_i(\varepsilon)$ and $c_j(\varepsilon)$.*

Proof. If $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} \neq 0$ then the equation

$$\left[\left(\zeta + \varepsilon X^{(1)}(\cdot) \right)^n \right]_{\mathbf{0}} = 0, \quad (2.8)$$

with $X^{(1)}$ defined in (2.4), admits at least one non-zero real solution ζ_1 of order ε . Thus, one can write

$$\left(\zeta + \varepsilon X^{(1)}(\cdot) \right)^n = \left(\zeta - \zeta_1 + \zeta_1 + \varepsilon X^{(1)}(\cdot) \right)^n = \sum_{k=0}^n \binom{n}{k} (\zeta - \zeta_1)^{n-k} \left(\zeta_1 + \varepsilon X^{(1)}(\cdot) \right)^k,$$

so that, by using that $[(\zeta_1 + \varepsilon X^{(1)}(\cdot))^n]_{\mathbf{0}} = 0$, one has

$$\left[\left(\zeta + \varepsilon X^{(1)}(\cdot) \right)^n \right]_{\mathbf{0}} = (\zeta - \zeta_1) \sum_{k=0}^{n-1} \binom{n}{k} (\zeta - \zeta_1)^{n-1-k} \left[\left(\zeta_1 + \varepsilon X^{(1)}(\cdot) \right)^k \right]_{\mathbf{0}}.$$

For $\zeta = \zeta_1$ one has

$$\sum_{k=0}^{n-1} \binom{n}{k} (\zeta - \zeta_1)^{n-1-k} \left[\left(\zeta_1 + \varepsilon X^{(1)}(\cdot) \right)^k \right]_{\mathbf{0}} = \binom{n}{n-1} \left[\left(\zeta_1 + \varepsilon X^{(1)}(\cdot) \right)^{n-1} \right]_{\mathbf{0}} > 0,$$

which shows that ζ_1 is a simple root of the equation (2.8), and hence also of the bifurcation equation (2.5). \blacksquare

The following result extends Lemma 2.2, and shows that the bifurcation equation admits a unique solution for any f .

Lemma 2.4. *The bifurcation equation (2.5) admits one and only one solution $c(\varepsilon)$. One has $c(\varepsilon) = \zeta_1 \varepsilon + O(\varepsilon^2)$, with $\zeta_1 \neq 0$ if $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} \neq 0$ and $c(\varepsilon) = O(\varepsilon^2)$ if $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} = 0$.*

Proof. By Lemma 2.2 we can confine ourselves to the case $[(X^{(1)}(\cdot))^n]_{\mathbf{0}} \neq 0$. In that case we know by Lemma 2.3 that all the solutions of the bifurcation equation are of order ε and separated by order ε . Hence we can set $\zeta = \zeta_1 \varepsilon + \zeta'$, with $\zeta' = o(\varepsilon)$, so that one has

$$\begin{aligned} \frac{\partial}{\partial c} F(\varepsilon, c_0 + \zeta) &= g_0 \frac{\partial}{\partial \zeta} \left[\left(\zeta + \varepsilon X^{(1)}(\cdot) \right)^n \right]_{\mathbf{0}} + O(\varepsilon^{n+1}) \\ &= g_0 \varepsilon^n \left[\left(\zeta_1 + X^{(1)}(\cdot) \right)^{n-1} \right]_{\mathbf{0}} + O(\varepsilon^{n+1}), \end{aligned}$$

which is strictly positive for $\varepsilon > 0$ small enough. Thus, at fixed ε , $F(\varepsilon, c)$ is increasing in c : this yields that the solution $c = c(\varepsilon)$ of $F(\varepsilon, c) = 0$ is unique for ε small enough. \blacksquare

3 Multiscale analysis and diagrammatic expansion

In this section we develop a diagrammatic representation for the quasi-periodic solution. There are a few differences with respect to [8], as we expand simultaneously both X and c in terms of the perturbation parameter ε . As in [8], the expansions we shall find are not power series expansions. See also [5, 4] for analogous situations; note, however, that in our case, as in [7], no fractional expansions arise.

We can summarise the results of the previous section as follows. The bifurcation equation (2.5) admits a unique solution $c(\varepsilon) = c_0 + \zeta(\varepsilon)$ such that $\zeta(0) = 0$. Let $X^{(1)}$ the zero-average function such that $\dot{X}^{(1)} = f$. If $[(X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} \neq 0$ then one has $\zeta(\varepsilon) = \zeta_1 \varepsilon + o(\varepsilon)$, where $\varepsilon \zeta_1$ is the unique solution of the equation (2.8).

To make notation uniform we can set in the following $\zeta_1 = 0$ if $[(X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} = 0$. By defining

$$P_{\mathbf{n}}(\zeta) = \left[\left(\zeta + X^{(1)}(\cdot) \right)^{\mathbf{n}} \right]_{\mathbf{0}}, \quad P'_{\mathbf{n}}(\zeta) = \frac{d}{d\zeta} P_{\mathbf{n}}(\zeta), \quad (3.1)$$

if $[(X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} \neq 0$ one has $P'_{\mathbf{n}}(\zeta_1) \neq 0$ by Lemma 2.3, and if $[(X^{(1)}(\cdot))^{\mathbf{n}}]_{\mathbf{0}} = 0$ one has $P'_{\mathbf{n}}(\zeta_1) = P'_{\mathbf{n}}(0) = \mathbf{n}[(X^{(1)}(\cdot))^{\mathbf{n}-1}]_{\mathbf{0}} \neq 0$ because \mathbf{n} is odd. Set $a = g_0 P'_{\mathbf{n}}(\zeta_1)$ in both cases; then $a \neq 0$.

In the following we give the diagrammatic rules in detail to make the exposition self-contained, and we only stress where the main differences lie with respect to the expansion of [8]. From a technical point of view, the diagrammatic analysis turns out a bit more involved, as it requires further expansions, and hence more labels to be assigned to the diagrams. On the other hand, eventually there is the advantage that one has a completely constructive algorithm to determine the solution within any fixed accuracy and an explicitly computable value for the maximal size of the allowed perturbation.

A graph is a connected set of points and lines. A *tree* θ is a graph with no cycle, such that all the lines are oriented toward a unique point (*root*) which has only one incident line (root line). All the points in a tree except the root are called *nodes*. The orientation of the lines in a tree induces a partial ordering relation (\preceq) between the nodes. Given two nodes v and w , we shall write $w \prec v$ every time v is along the path (of lines) which connects w to the root.

We call $E(\theta)$ the set of *end nodes* in θ , that is the nodes which have no entering line, and $V(\theta)$ the set of *internal nodes* in θ , that is the set of nodes which have at least one entering line. Set $N(\theta) = E(\theta) \amalg V(\theta)$. With each end node v we associate a *mode* label $\nu_v \in \mathbb{Z}^d$. For all $v \in N(\theta)$ denote with s_v the number of lines entering the node v ; for $v \in V(\theta)$ one has $s_v \geq \mathbf{n}$.

With respect to [8] the mode label of the end nodes can be $\mathbf{0}$. This only occurs when $\zeta_1 \neq 0$. Hence we define $E_0(\theta) = \{v \in E(\theta) : \nu_v = \mathbf{0}\}$ and $E_1(\theta) = \{v \in E(\theta) : \nu_v \neq \mathbf{0}\}$; we can set $E_0(\theta) = \emptyset$ if $\zeta_1 = 0$. Define also $L_0(\theta) = \{\ell \in L(\theta) : \ell \text{ exits a node } v \in E_0(\theta)\}$. If $|N(\theta)| = 1$ one requires $E_1(\theta) = N(\theta)$ and hence $E_0(\theta) = \emptyset$.

We denote with $L(\theta)$ the set of lines in θ . Since a line ℓ is uniquely identified with the node v which it leaves, we may write $\ell = \ell_v$. With each line ℓ we associate a *momentum* label $\nu_\ell \in \mathbb{Z}^d$ and a *scale* label $n_\ell \in \mathbb{Z}_+ \cup \{-1\}$. We set $n_\ell = -1$ when $\nu_\ell = \mathbf{0}$ (note that $\nu_\ell = \mathbf{0}$ was not allowed in [8]).

The modes of the end nodes and the momenta of the lines are related as follows: if $\ell = \ell_v$ one has

$$\boldsymbol{\nu}_\ell = \sum_{w \in E(\theta): w \preceq v} \boldsymbol{\nu}_w. \quad (3.2)$$

Given a tree θ we set $\Lambda_0(\theta) = \{\ell \in L(\theta) : \boldsymbol{\nu}_\ell = \mathbf{0}\} \setminus L_0(\theta) = \{\ell \in L(\theta) : \boldsymbol{\nu}_\ell = \mathbf{0} \text{ and } \ell \text{ does not exit an end node}\}$, and define the *order* of θ as

$$k(\theta) = |N(\theta)| - \mathbf{n} |\Lambda_0(\theta)|, \quad (3.3)$$

and the *total momentum* of θ as $\boldsymbol{\nu}(\theta) = \boldsymbol{\nu}_{\ell_0}$, if ℓ_0 is the root line of θ .

We call *equivalent* two trees which can be transformed into each other by continuously deforming the lines in such a way that they do not cross each other. Let $\mathcal{T}_{k,\boldsymbol{\nu}}$ be the set of inequivalent trees of order k and total momentum $\boldsymbol{\nu}$, that is the set of inequivalent trees θ such that $k(\theta) = k$ and $\boldsymbol{\nu}(\theta) = \boldsymbol{\nu}$.

A *cluster* T on scale n is a maximal set of nodes and lines connecting them such that all the lines have scales $n' \leq n$ and there is at least one line with scale n . The lines entering the cluster T and the possible line coming out from it (unique if existing at all) are called the external lines of the cluster T . Given a cluster T on scale n , we shall denote by $n_T = n$ the scale of the cluster. We call $V(T)$, $E(T)$, and $L(T)$ the set of internal nodes, of end nodes, and of lines of T , respectively; the external lines of T do not belong to $L(T)$.

We call *self-energy cluster* any cluster T such that T has only one entering line ℓ_T^2 and one exiting line ℓ_T^1 , and one has $\sum_{v \in E(T)} \boldsymbol{\nu}_v = \mathbf{0}$ (and hence $\boldsymbol{\nu}_{\ell_T^1} = \boldsymbol{\nu}_{\ell_T^2}$ by (3.2)). Set $x_T = \boldsymbol{\omega} \cdot \boldsymbol{\nu}_{\ell_T^2}$.

We call *excluded* a node v such that $\boldsymbol{\nu}_{\ell_v} = \mathbf{0}$, $s_v = \mathbf{n}$, at least $\mathbf{n} - 1$ lines entering v do exit end nodes, and the other line ℓ' entering v either also exits an end node or has momentum $\boldsymbol{\nu}_{\ell'} = \mathbf{0}$.

Let $\mathfrak{T}_{k,\boldsymbol{\nu}}$ be the set of *renormalised trees* in $\mathcal{T}_{k,\boldsymbol{\nu}}$, i.e. of trees in $\mathcal{T}_{k,\boldsymbol{\nu}}$ which contain neither any self-energy clusters nor any excluded nodes. Define also \mathfrak{R}_n as the set of renormalised self-energy clusters on scale n , i.e. of self-energy clusters on scale n which contain neither any further self-energy clusters nor any excluded nodes.

Lemma 3.1. *One has $k(\theta) \geq 1$ if $\boldsymbol{\nu}(\theta) \neq \mathbf{0}$ and $k(\theta) \geq 2$ if $\boldsymbol{\nu}(\theta) = \mathbf{0}$.*

Proof. By induction on $|N(\theta)|$. If $|N(\theta)| = 1$ then $N(\theta) = E_1(\theta)$ and hence $\Lambda_0(\theta) = \emptyset$, so that $k(\theta) = 1$. If $|N(\theta)| > 1$ let v_0 be the last node of θ , that is the node which the root line of θ exits, and let $\theta_1, \dots, \theta_s$ the trees with root in v_0 . Note that $s \geq \mathbf{n}$. Then one has $|N(\theta)| = 1 + |N(\theta_1)| + \dots + |N(\theta_s)|$, while $|\Lambda_0(\theta)| = |\Lambda_0(\theta_1)| + \dots + |\Lambda_0(\theta_s)|$ if $\boldsymbol{\nu}(\theta) \neq \mathbf{0}$ and $|\Lambda_0(\theta)| = 1 + |\Lambda_0(\theta_1)| + \dots + |\Lambda_0(\theta_s)|$ if $\boldsymbol{\nu}(\theta) = \mathbf{0}$. In the first case one has $k = k(\theta) = 1 + k(\theta_1) + \dots + k(\theta_s) \geq 1 + \mathbf{n}$. In the second case one has $k = k(\theta) = 1 + k(\theta_1) + \dots + k(\theta_s) - \mathbf{n} \geq 1 + s - \mathbf{n}$, so that $k \geq 1$. Moreover $k = 1$ would be possible only if $s = \mathbf{n}$ and $k(\theta_1) = \dots = k(\theta_s) = 1$. However, in such a case the lines entering v_0 would all exit end nodes and hence v_0 would be an excluded node. Thus, $k \geq 2$ if $\boldsymbol{\nu}(\theta) = \mathbf{0}$. ■

Lemma 3.2. *There exists a positive constant κ such that $|N(\theta)| \leq \kappa k(\theta)$ for any renormalised tree θ . One can take $\kappa = 3\mathbf{n}$.*

Proof. We give the proof for $\mathbf{n} > 1$ (the case $\mathbf{n} = 1$ being much easier; see [9]).

For $k(\theta) = 1$ the bound is trivially satisfied. We prove that for all $k \geq 2$, all $\nu \in \mathbb{Z}^d$ and all trees $\theta \in \mathfrak{T}_{k,\nu}$ one has $|N(\theta)| \leq ak(\theta) - b$, with $a = 3\mathbf{n}$ and $b = 4\mathbf{n}$.

The proof is by induction on k . For $k = 2$ it is just a check: if $\nu \neq \mathbf{0}$ one has $N(\theta) = 2$, while if $\nu = \mathbf{0}$ one has $N(\theta) = 2 + \mathbf{n} \leq 2\mathbf{n}$ for $\mathbf{n} > 1$.

By assuming that the bound holds for all $k < h$, one has for $\theta \in \mathfrak{T}_{h,\nu}$

$$|N(\theta)| = s_1 + \sum_{i=1}^{s_0} |N(\theta_i)|, \quad |\Lambda_0(\theta)| \leq 1 + \sum_{i=1}^{s_0} |\Lambda_0(\theta_i)|,$$

where $\ell_1, \dots, \ell_{s_0}$ are the lines in $\Lambda_0(\theta)$ closest to the root line of θ , $\theta_1, \dots, \theta_{s_0}$ are the trees with root lines $\ell_1, \dots, \ell_{s_0}$, respectively, and s_1 is the number of nodes which precede the root line of θ but not the root lines of $\theta_1, \dots, \theta_{s_0}$. By using the definition of order (3.3) one has

$$h = k(\theta) \geq s_1 - \mathbf{n} + \sum_{i=1}^{s_0} k(\theta_i).$$

Note that $k(\theta_i) \geq 2$ by Lemma 3.2, and hence $|N(\theta_i)| \leq ak(\theta_i) - b$ by the inductive hypothesis.

If $s_0 = 0$ then $\Lambda_0(\theta) \leq 1$, so that $|N(\theta)| \leq k(\theta) + \mathbf{n} = h + \mathbf{n} \leq 3\mathbf{n}h - 4\mathbf{n}$ for $\mathbf{n} > 1$ and $h > 1$.

If $s_0 \geq 1$ the inductive hypothesis yields

$$\begin{aligned} |N(\theta)| &\leq s_1 + a \sum_{i=1}^{s_0} k(\theta_i) - s_0 b \\ &\leq ak(\theta) - b - \left((a-1)s_1 + (s_0-1)b - a\mathbf{n} \right). \end{aligned}$$

Thus, the assertion follows if

$$(a-1)s_1 + (s_0-1)b - a\mathbf{n} = a(s_1 + s_0 - \mathbf{n}) + (b-a)(s_0-1) - s_1 - a \geq 0.$$

A key remark is that $s_0 + s_1 = \mathbf{n} + 1 + p$, with $p \geq 0$, because the root line of θ exits a node $v_0 \in V(\theta)$, so that $s_{v_0} \geq \mathbf{n}$. Then we can rewrite

$$a(s_1 + s_0 - \mathbf{n}) + (b-a)(s_0-1) - s_1 - a = ap - \mathbf{n} - p - 1 + s_0 + (b-a)(s_0-1).$$

If $p = 0$ then $s_1 + s_0 = \mathbf{n} + 1$; this yields $s_0 \geq 2$ (otherwise v_0 would be an excluded node), so that $ap - \mathbf{n} - p - 1 + s_0 + (b-a)(s_0-1) \geq 1 + (b-a) - \mathbf{n} \geq 1$. If $p \geq 1$ then $ap - \mathbf{n} - p - 1 + s_0 + (b-a)(s_0-1) \geq ap - \mathbf{n} - p = \mathbf{n}p + \mathbf{n}(p-1) + p(\mathbf{n}-1) \geq 2\mathbf{n} - 1$. \blacksquare

Let ψ be a non-decreasing C^∞ function defined in \mathbb{R}_+ , such that

$$\psi(u) = \begin{cases} 1, & \text{for } u \geq 1, \\ 0, & \text{for } u \leq 1/2, \end{cases} \quad (3.4)$$

and set $\chi(u) := 1 - \psi(u)$. For all $n \in \mathbb{Z}_+$ define $\chi_n(u) := \chi(u/4\alpha_n(\boldsymbol{\omega}))$ and $\psi_n(u) := \psi(u/4\alpha_n(\boldsymbol{\omega}))$, and set $\Xi_0(x) = \chi_0(|x|)$, $\Psi_0(x) = \psi_0(|x|)$, and

$$\Xi_n(x) = \chi_0(|x|) \cdots \chi_{n-1}(|x|) \chi_n(|x|), \quad \Psi_n(x) = \chi_0(|x|) \cdots \chi_{n-1}(|x|) \psi_n(|x|), \quad (3.5)$$

for $n \geq 1$. Then we define the node factor as

$$F_v = \begin{cases} -\frac{1}{s_v!} \frac{d^{s_v}}{dx^{s_v}} g(c_0), & v \in V(\theta), \\ f_{\boldsymbol{\nu}_v}, & v \in E_1(\theta), \\ \zeta_1, & v \in E_0(\theta), \end{cases} \quad (3.6)$$

and the propagator as

$$G_\ell = \begin{cases} G^{[n_\ell]}(\boldsymbol{\omega} \cdot \boldsymbol{\nu}_\ell; \varepsilon, c_0), & \boldsymbol{\nu}_\ell \neq \mathbf{0}, \\ -\frac{1}{a}, & \boldsymbol{\nu}_\ell = \mathbf{0}, \quad \ell \in \Lambda_0(\theta), \\ 1, & \boldsymbol{\nu}_\ell = \mathbf{0}, \quad \ell \in L_0(\theta), \end{cases} \quad (3.7)$$

with ζ_1 and a defined before and after (3.1), respectively, and $G^{[n]}(x; \varepsilon, c_0)$ recursively defined for $n \geq 0$ as

$$G^{[n]}(x; \varepsilon, c) = \frac{\Psi_n(x)}{ix(1 + i\varepsilon x) - \mathcal{M}^{[n-1]}(x; \varepsilon, c)}, \quad (3.8a)$$

$$\mathcal{M}^{[n]}(x; \varepsilon, c) = \mathcal{M}^{[n-1]}(x; \varepsilon, c) + \Xi_n(x) M^{[n]}(x; \varepsilon, c), \quad M^{[n]}(x; \varepsilon, c) = \sum_{T \in \mathfrak{R}_n} \text{Val}(T, x; \varepsilon, c), \quad (3.8b)$$

where $\mathcal{M}^{[-1]}(x; \varepsilon, c_0) = \varepsilon g_1(c_0) = 0$ and

$$\text{Val}(T, x; \varepsilon, c_0) = \left(\prod_{\ell \in L(T)} G_\ell \right) \left(\prod_{v \in N(T)} F_v \right) \quad (3.9)$$

is called the value of the self-energy cluster T .

Note that, with respect to the analogous formulae (3.5) and (3.6) of [8], both F_v and G_ℓ are computed at $c = c_0$. We prefer writing explicitly the dependence on c_0 (even if not necessary, as c_0 is fixed once and for all as in Assumption 2) simply in order to use the same notations as in [8].

Set

$$X_{\mathcal{V}}^{[k]} = \sum_{\theta \in \mathfrak{I}_{k, \boldsymbol{\nu}}} \text{Val}(\theta; \varepsilon, c_0), \quad \boldsymbol{\nu} \neq \mathbf{0}, \quad k \geq 1, \quad (3.10)$$

and

$$\zeta^{[k]} = \sum_{\theta \in \mathfrak{I}_{k+n, \mathbf{0}}} \text{Val}(\theta; \varepsilon, c_0), \quad k \geq 2, \quad (3.11)$$

where the tree value $\text{Val}(\theta; \varepsilon, c_0)$ is defined as

$$\text{Val}(\theta, x; \varepsilon, c_0) = \left(\prod_{\ell \in L(\theta)} G_\ell \right) \left(\prod_{v \in N(\theta)} F_v \right). \quad (3.12)$$

Call $\mathfrak{N}_n(\theta)$ the number of lines $\ell \in L(\theta)$ such that $n_\ell \geq n$, and $\mathfrak{N}_n(T)$ the number of lines $\ell \in L(T)$ such that $n_\ell \geq n$, and set

$$M(\theta) = \sum_{v \in E(\theta)} |\nu_v|, \quad M(T) = \sum_{v \in E(T)} |\nu_v|. \quad (3.13)$$

Define the renormalised series

$$\bar{X}(\psi; \varepsilon) = \sum_{\nu \in \mathbb{Z}_+^d} e^{i\nu \cdot \psi} \bar{X}_\nu, \quad \bar{X}_\nu = \sum_{k=1}^{\infty} \varepsilon^k X_\nu^{[k]}. \quad (3.14)$$

and

$$\bar{\zeta}(\varepsilon) = \varepsilon \zeta_1 + \sum_{k=2}^{\infty} \varepsilon^k \zeta^{[k]}. \quad (3.15)$$

In the next section we shall prove first that the series (3.14) and (3.15) converge, then that the function $x_0(t) = c_0 + \bar{\zeta}(\varepsilon) + \bar{X}(\omega t; \varepsilon)$ solves the equations

$$i\omega \cdot \nu (1 + i\varepsilon \omega \cdot \nu) X_\nu + \varepsilon [g(c_0 + \zeta(\varepsilon) + X(\cdot; \varepsilon))]_\nu = \varepsilon f_\nu, \quad \nu \neq \mathbf{0}, \quad (3.16a)$$

$$[g(c_0 + \zeta(\varepsilon) + X(\cdot; \varepsilon))]_{\mathbf{0}} - f_{\mathbf{0}} = 0, \quad (3.16b)$$

and hence the equation (1.1).

4 Bounds

The proof of the convergence of the renormalised series proceeds as in [8]. We confine ourselves to state the basic steps of the proof, without giving the details, except when the discussion departs from [8].

Lemma 4.1. *For any renormalised tree θ , one has $\mathfrak{N}_n(\theta) \leq 2^{-(n-2)} M(\theta)$.*

Proof. The same as the proof of Lemma 3.1 in [8]. ■

Lemma 4.2. *Assume there exists a constant C_0 such that $|G^{[n]}(x; \varepsilon, c_0)| \leq C_0/\alpha_n(\omega)$ for all $n \in \mathbb{Z}_+$. Then there exists $\varepsilon_0 > 0$ such that, for all all $|\varepsilon| < \varepsilon_0$, the series (3.14) and (3.15) converge. Moreover the series (3.14) is analytic in ψ .*

Proof. Let θ be a tree in $\mathfrak{T}_{k,\nu}$ and set $K = |N(\theta)|$. One reasons as in [8] to prove that

$$|\text{Val}(\theta; \varepsilon, c)| \leq \tilde{C}_0^K D_0^K \alpha_{n_0}^{-K}(\omega) e^{-\xi M(\theta)/2},$$

where $D_0 = \max\{\Gamma, \Phi\}$, $\tilde{C}_0 = \max\{C_0, a\}$, and $n_0 \in \mathbb{N}$ such that

$$4 \sum_{n=n_0+1}^{\infty} \frac{1}{2^n} \log \frac{1}{\alpha_n(\omega)} \leq \frac{\xi}{2}.$$

By Lemma 3.2 one has $K \leq \kappa k$. The number of trees of order k with fixed mode labels is hence bounded by C_2^k for a suitable constant C_2 , and the sum over the mode labels can be performed by using half the exponent in the decaying factor $e^{-\xi M(\theta)/2}$. This gives

$$\left| X_{\mathcal{V}}^{[k]} \right| \leq C^k e^{-\xi |\nu|/4},$$

for some constant C , which yields the convergence of the series (3.14) to a function analytic in ψ . In the same way one obtains

$$\left| \zeta^{[k]} \right| \leq C^k$$

with the same constant C , and this yields the convergence of the series (3.15). \blacksquare

Lemma 4.3. *For any self-energy cluster $T \in \mathfrak{R}_n$ such that $\Xi_n(x_T) \neq 0$, one has $M(T) \geq 2^{n-1}$ and $\mathfrak{N}_p(T) \leq 2^{-(p-2)} M(T)$ for all $p \leq n$.*

Proof. The same as the proof of Lemma 3.3 of [8]. \blacksquare

Lemma 4.4. *Assume the propagators $G^{[p]}(x; \varepsilon, c_0)$ are differentiable in x and there exist constants C_0 and C_1 such that $|G^{[p]}(x; \varepsilon, c_0)| \leq C_0/\alpha_p(\omega)$ and $|\partial_x G^{[p]}(x; \varepsilon, c_0)| \leq C_1/\alpha_p^3(\omega)$ for all $p < n$. Then there exists $\varepsilon_0 > 0$ such that, for all $|\varepsilon| < \varepsilon_0$, the function $x \mapsto M^{[n]}(x; \varepsilon, c_0)$ is differentiable, and one has*

$$\left| M^{[n]}(x; \varepsilon, c_0) \right|, \left| \partial_x M^{[n]}(x; \varepsilon, c_0) \right| \leq D_1 |\varepsilon|^2 e^{-D_2 2^n},$$

for some positive constants D_1 and D_2 .

Proof. Proceed as in the proof of Lemma 3.4 of [8], by taking into account the differences already pointed out in the proof of Lemma 4.2. \blacksquare

Lemma 4.5. *Assume there exists a constant C_0 such that $|G^{[p]}(x; \varepsilon, c_0)| \leq C_0/\alpha_p(\omega)$ for all $p < n$. Then one has $(\mathcal{M}^{[p]}(x; \varepsilon, c_0))^* = \mathcal{M}^{[p]}(-x; \varepsilon, c_0)$ for all $p \leq n$.*

Proof. As the proof of Lemma 3.5 of [8]. \blacksquare

Lemma 4.6. *Let ε_0 as in Lemma 4.1. For all $n \in \mathbb{Z}_+$ the function $x \mapsto \mathcal{M}^{[n]}(x; \varepsilon, c_0)$ is differentiable and one has $|ix(1 + i\varepsilon x) - \mathcal{M}^{[n]}(x; \varepsilon, c_0)| \geq |x|/2$ for all $|\varepsilon| < \varepsilon_0$.*

Proof. As the proof of Lemma 3.6 of [8]. ■

By collecting together the results above we can prove that the series (3.14) and (3.15) converge. The proof proceeds as follows. Convergence holds if the hypotheses of Lemma 4.2 are verified. So we have to prove that the propagators satisfy the bounds $|G^{[n]}(x; \varepsilon, c_0)| \leq C_0/\alpha_n(\omega)$. More precisely, we can prove by induction on n that the propagators satisfy the bounds

$$\left| G^{[n]}(x; \varepsilon, c_0) \right| \leq \frac{C_0}{\alpha_n(\omega)}, \quad \left| \partial_x G^{[n]}(x; \varepsilon, c_0) \right| \leq \frac{C_1}{\alpha_n^3(\omega)}, \quad (4.1)$$

for suitable constants C_0 and C_1 . For $n = 0$ the check is trivial. Then we assume that the bounds are satisfied up to scale $n - 1$, and we write $G^{[n]}(x; \varepsilon, c_0)$ according to (3.8). Now, $G^{[n]}(x; \varepsilon, c_0)$ depends on the quantities $\mathcal{M}^{[p]}(x; \varepsilon, c_0)$ with $p < n$, which in turn depend on the propagators $G^{[p']}(x; \varepsilon, c_0)$ with $p' \leq p$. Thus, by the inductive hypothesis we can apply Lemma 4.4 and Lemma 4.5 to deduce Lemma 4.6. This implies that also $G^{[n]}(x; \varepsilon, c_0)$ satisfies the bounds (4.1). Smoothness in ε of the series is discussed in the same way, proving by induction that the derivatives of the propagators satisfy the bounds $|\partial_\varepsilon^m G^{[n]}(x; \varepsilon, c_0)| \leq K_m/\alpha_n^m(\omega)$ for all $m \geq 1$ and for suitable constants K_m (we refer to [8] for further details).

Finally, by proceeding as in the proof of Lemma 3.7 of [8], we can prove that the function $x_0(t) = c_0 + \zeta(\varepsilon) + X(\omega t; \varepsilon)$ solves the equations (3.16). As far as equation (3.16a) is concerned the proof proceeds as in the proof of Lemma 3.7 of [8]. To deal with equation (3.16b), we write $[g(c_0 + \zeta(\varepsilon) + X(\cdot; \varepsilon))]_0$ according to (2.6) and expand $X(\cdot; \varepsilon) = \overline{X}(\cdot; \varepsilon)$ according to (3.14) and (3.10). Then we obtain

$$g_0 P'_n(\zeta_1) \left(\zeta_k + \sum_{\theta \in \mathfrak{I}_{k+n,0}} \text{Val}(\theta; \varepsilon, c_0), \right) \quad k \geq 2, \quad (4.2)$$

which is identically zero if ζ_k is defined by (3.11).

References

- [1] M.V. Bartuccelli, J.H.B. Deane, G. Gentile, *Globally and locally attractive solutions for quasi-periodically forced systems*, J. Math. Anal. Appl. **328** (2007), no. 1, 699–714.
- [2] E. Brieskorn, H. Knörrer, *Plane algebraic curves*, Birkhäuser, Basel, 1986.
- [3] L.O. Chua, A. Ushida, *Algorithms for computing almost periodic steady-state response of nonlinear systems to multiple input frequencies*, IEEE Trans. Circuits and Systems **28** (1981), no. 10, 953–971.
- [4] L. Corsi, G. Gentile, *Melnikov theory to all orders and Puiseux series for subharmonic solutions*, J. Math. Phys. **49** (2008), no. 11, 112701, 29 pp.
- [5] G. Gallavotti, G. Gentile, A. Giuliani, *Fractional Lindstedt series*, J. Math. Phys. **47** (2006), no. 1, 012702, 33 pp.
- [6] G. Gentile, *Quasi-periodic solutions for two-level systems*, Comm. Math. Phys. **242** (2003), no. 1-2, 221–250.

- [7] G. Gentile, D.A. Cortez, J.C.A. Barata, *Stability for quasi-periodically perturbed Hill's equations*, Comm. Math. Phys. **260** (2005), no. 2, 403–443.
- [8] G. Gentile, *Quasi-periodic motions in strongly dissipative forced systems*, Ergodic Theory Dynam. Systems, to appear.
- [9] G. Gentile, M.V. Bartuccelli, J.H.B. Deane, *Summation of divergent series and Borel summability for strongly dissipative differential equations with periodic or quasiperiodic forcing terms*, J. Math. Phys. **46** (2005), no. 6, 062704, 21 pp.
- [10] G. Gentile, M.V. Bartuccelli, J.H.B. Deane, *Quasiperiodic attractors, Borel summability and the Bryuno condition for strongly dissipative systems*, J. Math. Phys. **47** (2006), no. 7, 072702, 10 pp.
- [11] Ch. Kaas-Petersen, *Computation of quasiperiodic solutions of forced dissipative systems*, J. Comput. Phys. **58** (1985), no. 3, 395–408.
- [12] Ch. Kaas-Petersen, *Computation of quasiperiodic solutions of forced dissipative systems. II*, J. Comput. Phys. **64** (1986), no. 3, 433–442.