

Action minimizing properties and distances on the group of Hamiltonian diffeomorphisms

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February 20, 2010

In this article we prove that for a smooth fiberwise convex Hamiltonian, the asymptotic Hofer distance from the identity gives a strict upper bound to the value at 0 of Mather's β function, thus providing a negative answer to a question asked by K. Siburg in [15]. However, we show that equality holds if one considers the asymptotic distance defined in [17].

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[†]Supported also by *Fondation des Sciences Mathématiques de Paris*, *Herchel-Smith foundation* and *Agence Nationale de la Recherche* project "Hamilton-Jacobi et théorie KAM faible".

[‡]Supported also by *Agence Nationale de la Recherche* projects "Symplexe" and "Floer Power".

1 Introduction

The relationship between Aubry-Mather theory and the new tools of symplectic topology has attracted quite a bit of attention over the last years. These two approaches correspond to two different ways of looking at Hamiltonian systems. While the former investigates the dynamics of the system in the phase space, the latter takes a more global look at the topology of the path that the Hamiltonian flow describes in the group of Hamiltonian diffeomorphisms. Trying to relate and combine these “internal” and “external” information is a very intriguing task.

In this article we shall concentrate on the relation between the action minimizing properties of the flow of a convex Hamiltonian and its “asymptotic distance” from the identity.

Given a Tonelli Hamiltonian H (see section 3 for a precise definition) and its corresponding Lagrangian L (obtained by Legendre duality), Aubry-Mather theory associates the so-called β -function (or *effective Lagrangian*). Roughly speaking, the value of such a function represents the minimal average Lagrangian action needed to carry out motions with a prescribed “rotation vector”. The Legendre dual of the β -function is what is called the α -function (or *effective Hamiltonian*). See section 3 for precise definitions.

On the other hand the corresponding Hamiltonian flow φ_H^t determines a curve in the group of Hamiltonian diffeomorphisms. Recall that there are several natural metrics that can be defined on this group, in particular the so-called *Hofer distance* [10] and γ -distance¹ [17] (see section 2).

Studying the connection between these objects arises quite naturally and has indeed been studied by several authors in the last years.

One question, for instance, concerns the relationship between $\beta_H(0)$, associated to a Tonelli Hamiltonian H , and the asymptotic Hofer distance from the identity of its time-one flow map φ_H^1 . Observe that the definitions need to be adjusted since φ_H^1 is not compactly supported, while the Hofer distance is only defined for compactly supported Hamiltonians.

In [15] K. Siburg proves, in the case of Hamiltonian diffeomorphisms on the cotangent disc bundle generated by a convex Hamiltonian H , that the asymptotic Hofer distance yields an upper bound for $\beta_H(0)$ (see Proposition 2) and asks whether or not equality holds. In this paper we show that Siburg’s question has a negative answer (Corollary 2), by constructing examples of convex Hamiltonian diffeomorphisms for which the asymptotic Hofer distance from the identity is

¹Sometimes called *Viterbo distance*.

strictly greater than the asymptotic γ -distance (Theorem 2). However, Siburg's question has a positive answer, provided the asymptotic Hofer distance is replaced by the asymptotic γ -distance (Proposition 2). Moreover, we extend these results to the case of general Hamiltonian diffeomorphisms generated by autonomous Tonelli Hamiltonians (Theorem 1 and Corollary 1).

Observe that Corollary 2 is also stated in [7]. However, as the author kindly confirmed to us, there is a gap in the proof of [7, Proposition 7]. Although our proof goes along completely different lines, we are grateful to X. Cui for drawing our attention to this problem.

2 Metric structures on the group of Hamiltonian diffeomorphisms

We first define the group of Hamiltonian diffeomorphisms and two metrics on this group: *Hofer distance* and *γ -distance*.

Let $\|\cdot\|$ denote the standard metric on the n -dimensional torus $\mathbb{T}^n \simeq \mathbb{R}^n/\mathbb{Z}^n$ and $\omega = -d\lambda$ the canonical symplectic structure on $\mathbb{T}^*\mathbb{T}^n$, where $\lambda = \sum_{j=1}^n p_j dq_j$ is the Liouville form on $\mathbb{T}^*\mathbb{T}^n$. We denote by \mathcal{H}_0 the set of *admissible* time-dependent Hamiltonians $H \in C^2(\mathbb{T}^*\mathbb{T}^n \times \mathbb{T})$, such that $H_t(q, p) := H(q, p, t)$ has compact support. For each $H \in \mathcal{H}_0$ we consider the corresponding Hamiltonian flow φ_H^t and denote by $\varphi_H := \varphi_H^1$ its time-one map. The *group of Hamiltonian diffeomorphisms* $\mathcal{Ham}_0(\mathbb{T}^*\mathbb{T}^n) := \mathcal{Ham}_0(\mathbb{T}^*\mathbb{T}^n, \omega)$ is the set of all Hamiltonian diffeomorphisms $\varphi : \mathbb{T}^*\mathbb{T}^n \rightarrow \mathbb{T}^*\mathbb{T}^n$ that are obtained as time-one maps of elements in \mathcal{H}_0 , *i.e.* $\varphi = \varphi_H$ for some $H \in \mathcal{H}_0$.

We shall now define the Hofer and γ -distances for elements in $\mathcal{Ham}_0(\mathbb{T}^*\mathbb{T}^n)$.

2.1 The Hofer distance

This first metric structure on the group of compactly supported Hamiltonian diffeomorphisms was defined by Hofer (see [10]). Consider a path in the group of compactly supported Hamiltonian diffeomorphisms, given by an admissible Hamiltonian H . One first defines the length $\ell(H)$ of this Hamiltonian path, by setting

$$\ell(H) := \int_{\mathbb{T}} \text{Osc } H_t dt,$$

where $\text{Osc } H_t := \max_{\mathbb{T}^*\mathbb{T}^n} H_t - \min_{\mathbb{T}^*\mathbb{T}^n} H_t$ denotes the oscillation of H_t .

Definition 1 (Hofer distance). *The Hofer distance from the identity (or energy) of an element $\varphi \in \mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$ is given by*

$$d^H(\text{id}, \varphi) := \inf \{ \ell(H) : H \in \mathcal{H}_0 \text{ and } \varphi = \varphi_H \}.$$

This extends to a distance on $\mathcal{H}am_0(\mathbb{T}^\mathbb{T}^n)$: if $\varphi, \psi \in \mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$, then $d^H(\varphi, \psi) := d^H(\text{id}, \psi \circ \varphi^{-1})$.*

It is easy to verify that if $H, K \in \mathcal{H}_0$, then

$$d^H(\varphi_H^1, \varphi_K^1) \leq \|H - K\|_{C^0}.$$

Observe that this definition only considers the flow of a given Hamiltonian at time $t = 1$. In the study of the dynamics, however, one is interested in the long time behaviour of the system and it would be more relevant to get global information, such as, for instance, the *asymptotic Hofer distance from the identity* introduced by Bialy and Polterovich in [2].

Definition 2 (Asymptotic Hofer distance). *Let φ be a Hamiltonian diffeomorphism in $\mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$. The asymptotic Hofer distance from the identity is:*

$$d_\infty^H(\text{id}, \varphi) := \lim_{k \rightarrow +\infty} \frac{d^H(\text{id}, \varphi^k)}{k}.$$

It follows from the triangle inequality that the above limit exists and that $d_\infty^H \leq d^H$.

2.2 The γ -distance

One can also introduce another metric on $\mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$, commonly referred to as γ -distance. First of all, let us recall the following construction (see [17] for more details). Let \mathcal{L}_0 denote the set of Lagrangian submanifolds Λ of $\mathbb{T}^*\mathbb{T}^n$, which are Hamiltonianly isotopic to the zero section $\mathcal{O}_{\mathbb{T}^n}$, *i.e.* there exists a Hamiltonian isotopy φ^t such that $\Lambda = \varphi^1(\mathcal{O}_{\mathbb{T}^n})$. Consider $\Lambda \in \mathcal{L}_0$ and let $S_\Lambda : \mathbb{T}^n \times \mathbb{R}^k \rightarrow \mathbb{R}$ be a *generating function quadratic at infinity* (GFQI) for Λ (see [17] for the definition). Since Λ is Hamiltonianly isotopic to $\mathcal{O}_{\mathbb{T}^n}$, then S_Λ is unique up to some elementary operations. Moreover, if we denote by $S_\Lambda^\lambda := \{(q; \xi) \in \mathbb{T}^n \times \mathbb{R}^k : S_\Lambda(q; \xi) \leq \lambda\}$, then for sufficiently large $c \in \mathbb{R}$ we have that $H^*(S_\Lambda^c, S_\Lambda^{-c}) \simeq H^*(\mathbb{T}^n) \otimes H^*(D^-, \partial D^-)$, where D^- is the unit disc of the negative eigenspace of the quadratic form B associated to S_Λ . Therefore to each cohomology class

$\alpha \in H^*(\mathbb{T}^n) \setminus \{0\}$, one can associate the image $\alpha \otimes T$ (T is a chosen generator of $H^*(D^-, \partial D^-) \simeq \mathbb{Z}$) and, by min-max methods, a critical level $c(\alpha, S_\Lambda)$ (see [17, p. 690-693] for more details).

Let us now consider $\varphi \in \mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$. Recall that its graph

$$\Gamma(\varphi) = \{(z, \varphi(z)) : z \in \mathbb{T}^*\mathbb{T}^n\}$$

is a Lagrangian submanifold of $\mathbb{T}^*\mathbb{T}^n \times \overline{\mathbb{T}^*\mathbb{T}^n}$, where $\overline{\mathbb{T}^*\mathbb{T}^n}$ denotes $\mathbb{T}^*\mathbb{T}^n$ with the symplectic form $-\omega$ (see for example [5]). Since $\mathbb{T}^*\mathbb{T}^n \times \overline{\mathbb{T}^*\mathbb{T}^n}$ is covered by $\mathbb{T}^*(\Delta_{\mathbb{T}^*\mathbb{T}^n})$, where $\Delta_{\mathbb{T}^*\mathbb{T}^n}$ is the diagonal, we may lift $\Gamma(\varphi)$ to $\tilde{\Gamma}(\varphi)$, which is still a Lagrangian submanifold in $\mathbb{T}^*(\Delta_{\mathbb{T}^*\mathbb{T}^n})$. Moreover, since φ has compact support, we can compactify both $\tilde{\Gamma}(\varphi)$ and $\Delta_{\mathbb{T}^*\mathbb{T}^n}$ and we obtain a Lagrangian submanifold $\bar{\Gamma}(\varphi)$ in $\mathbb{T}^*(S^n \times \mathbb{T}^n)$. In [17, p. 697 and p. 706] the second author defined the following distance from the identity.

Definition 3 (γ -distance). *Let $\varphi \in \mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$. The γ -distance of φ from the identity is given by:*

$$\gamma(\text{id}, \varphi) := c(\mu_{\mathbb{T}^n} \otimes \mu_{S^n}, \bar{\Gamma}(\varphi)) - c(1 \otimes 1, \bar{\Gamma}(\varphi)).$$

In particular, this can be easily extended to a distance on the all group: if $\varphi, \psi \in \mathcal{H}am_0(\mathbb{T}^\mathbb{T}^n)$, then $\gamma(\varphi, \psi) := \gamma(\text{id}, \psi \circ \varphi^{-1})$.*

We also define

$$c_+(\varphi) := c(\mu_{\mathbb{T}^n} \otimes \mu_{S^n}, \bar{\Gamma}(\varphi))$$

and

$$c_-(\varphi) := c(1 \otimes 1, \bar{\Gamma}(\varphi))$$

so that $\gamma(\text{id}, \varphi) = c_+(\varphi) - c_-(\varphi)$.

It is possible to show again that $\gamma(\varphi_H^1, \varphi_K^1) \leq \|H - K\|_{C^0}$ for all $H, K \in \mathcal{H}_0$, and more precisely that $\gamma(\varphi_H^1, \varphi_K^1) \leq d^H(\varphi_H^1, \varphi_K^1)$ (see Proposition 1).

Analogously to what we have already seen for Hofer distance, one can introduce the *asymptotic γ -distance from the identity*:

Definition 4 (Asymptotic γ -distance). *Let φ be a Hamiltonian diffeomorphism in $\mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$. The asymptotic γ -distance from the identity is:*

$$\gamma_\infty(\text{id}, \varphi) := \lim_{k \rightarrow +\infty} \frac{\gamma(\text{id}, \varphi^k)}{k}.$$

Similarly,

$$c_{\pm, \infty}(\varphi) = \lim_{k \rightarrow +\infty} \frac{c_\pm(\varphi^k)}{k}.$$

Proposition 1. *For all $\varphi \in \mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$, $\gamma(\text{id}, \varphi) \leq d^H(\text{id}, \varphi)$. In particular, $\gamma_\infty(\text{id}, \varphi) \leq d_\infty^H(\text{id}, \varphi)$.*

This is an immediate consequence of [17, Proposition 4.6] and it was explicitly stated for example in [18, Proposition 2.15] or in [11, Proposition 1.52].

3 Action minimizing properties of convex Hamiltonians

In this section we want to recall some notions of Mather's theory of minimal action. In order to do this, we need to restrict our analysis to a special class of Hamiltonians $H : \mathbb{T}^*\mathbb{T}^n \times \mathbb{T} \rightarrow \mathbb{R}$, which are C^2 , strictly convex and superlinear in the fibers and have complete flows. These Hamiltonians, which are also called *Tonelli Hamiltonians*, play an important role in the study of classical mechanics and provide the setting in which *Mather's theory* and *Fathi's Weak KAM theory* have been developed (see for instance [14, 9]).

Let H be a Tonelli Hamiltonian and consider the associated Lagrangian $L : \mathbb{T}\mathbb{T}^n \times \mathbb{T} \rightarrow \mathbb{R}$, which is defined by Legendre duality by the formula:

$$L(q, v, t) = \sup\{\langle p, v \rangle - H(q, p, t) \mid p \in \mathbb{R}^n\}.$$

Recall that the associated Euler-Lagrange flow φ_L^t is obtained as solution of the equation $\frac{d}{dt} \frac{\partial L}{\partial v}(q, v, t) = \frac{\partial L}{\partial q}(q, v, t)$ and it is conjugated, via the Legendre transform $(q, v) \mapsto (q, \frac{\partial L}{\partial v}(q, v, t))$, to the Hamiltonian flow φ_H^t . Let μ be a probability measure on $\mathbb{T}\mathbb{T}^n \times \mathbb{T}$, which is invariant under the Euler-Lagrange flow (i.e. $(\varphi_L^t)_*\mu = \mu$). We define its *average action* as

$$A_L(\mu) := \int_{\mathbb{T}\mathbb{T}^n \times \mathbb{T}} L(q, v, t) d\mu.$$

Let us denote by $\mathfrak{M}(L)$ the space of invariant probability measures on $\mathbb{T}\mathbb{T}^n \times \mathbb{T}$, with finite average action. Given any $\mu \in \mathfrak{M}(L)$ we can define its *rotation vector* or *Schwartzman's asymptotic cycle* as the unique $\rho(\mu) \in H_1(\mathbb{T}^n; \mathbb{R})$ that satisfies

$$\int_{\mathbb{T}\mathbb{T}^n \times \mathbb{T}} \eta(q, t) \cdot (v, 1) d\mu = \langle \rho(\mu), [\eta]_{\mathbb{T}^n} \rangle + [\eta]_{\mathbb{T}}$$

for any closed 1-form η on $\mathbb{T}^n \times \mathbb{T}$, where $[\eta] = ([\eta]_{\mathbb{T}^n}, [\eta]_{\mathbb{T}}) \in H_1(\mathbb{T}^n \times \mathbb{T}; \mathbb{R}) \simeq H_1(\mathbb{T}^n; \mathbb{R}) \times \mathbb{R}$ is the de-Rham cohomology of η . It is possible to show [14] that

the map $\rho : \mathfrak{M}(L) \longrightarrow H_1(\mathbb{T}^n; \mathbb{R})$ is surjective and hence there exist invariant probability measures for each rotation vector. Let us consider the minimal value of the average action A_L over the set of probability measures with a given rotation vector. This minimum exists because of the lower semicontinuity of the action functional on the set $\mathfrak{M}(L)$ for the weak-* topology (see [14]):

$$\begin{aligned} \beta_H : H_1(\mathbb{T}^n; \mathbb{R}) &\longrightarrow \mathbb{R} \\ h &\longmapsto \min_{\mu \in \mathfrak{M}(L) : \rho(\mu) = h} A_L(\mu). \end{aligned}$$

This function β_H is what is generally known as β -function or *effective Lagrangian*. A measure $\mu \in \mathfrak{M}(L)$ realizing such a minimum amongst all invariant probability measures with the same rotation vector, *i.e.* $A_L(\mu) = \beta(\rho(\mu))$, is called an *action minimizing measure with rotation vector* $\rho(\mu)$. The β -function is convex and therefore one can consider its *conjugate* function (given by Fenchel duality) $\alpha_H : H^1(\mathbb{T}^n; \mathbb{R}) \longrightarrow \mathbb{R}$ defined by

$$\alpha_H(c) := \max_{h \in H_1(\mathbb{T}^n; \mathbb{R})} (\langle c, h \rangle - \beta_H(h)).$$

This function is generally called α -function or *effective Hamiltonian*. See [14] for more details.

Remark 1. It turns out that α_H coincides with the symplectic homogenization of H , usually denoted \overline{H} (see [19, Proposition 9.3]).

More precisely, in the autonomous case, we just have $\overline{H}(p) = \alpha_H(p)$ while in the non-autonomous case, we reduce ourselves to the autonomous one by setting $K(q, p, t, \tau) = \tau + H(q, p, t)$. Note that $\overline{K}(p, \tau)$ is well defined and equal to $\tau + \overline{H}(p)$ for some function \overline{H} . For this function, we have again $\overline{H}(p) = \alpha_H(p)$.

4 Main results

In this section we want to study the connection between the Hofer and γ -distance on one hand, and Mather's theory of minimal action on the other hand. We shall then state our main results. Let us start by observing that Tonelli Hamiltonians clearly do not belong to \mathcal{H}_0 , since they lack compact support, and hence their time-one maps are not element of $\mathcal{H}am_0(\mathbb{T}^*\mathbb{T}^n)$. Therefore, we need to restrict them to compact subsets of $\mathbb{T}^*\mathbb{T}^n$ and, as done by Siburg in [15], consider "nice" compactly supported extensions, such that the Hofer and γ -metrics are independent of the choice of the extension. In the autonomous case this can be achieved

by using the conservation of energy to obtain well-defined “truncated” flows and then smooth them out. We shall then see how the same proof extends to the setting considered in [15].

Let us consider $H(q, p)$ an autonomous Tonelli Hamiltonian on $T^*\mathbb{T}^n$. Then, for each r and each sufficiently small $\varepsilon > 0$, let us consider functions $f_{r,\varepsilon}$ such that:

- $f_{r,\varepsilon}(s) = s$ for $s \leq r$;
- $f_{r,\varepsilon}(s) \leq r + \varepsilon$ for all $s \in \mathbb{R}$ and $f_{r,\varepsilon}(s) \equiv r$ for $s \geq r + \varepsilon$;
- $\sqrt{\varepsilon}|f'_{r,\varepsilon}(s)| \leq 1$ for all $s \in \mathbb{R}$ (this assumption will only be used in section 4.3 to define the Calabi invariant).

Now, we can define new Hamiltonians, given by $H_{r,\varepsilon} = f_{r,\varepsilon}(H)$. If we denote by $S_r := \{(q, p) \in T^*\mathbb{T}^n : H(q, p) \leq r\}$, then our new Hamiltonians will satisfy the following conditions:

- $H_{r,\varepsilon} - r$ is supported in $S_{r+\varepsilon}$;
- $H_{r,\varepsilon}$ coincides with H on S_r ;
- $H_{r,\varepsilon}$ is bounded everywhere by $r + \varepsilon$ and satisfies the condition $\sqrt{\varepsilon}\|d_p H_{r,\varepsilon}\| \leq C$, where $C = \sup_{S_{r+\varepsilon}} \|d_p H\|$.

We shall denote the set of all these possible extensions by $\mathcal{H}_{r,\varepsilon}(H)$. We can now define Hofer and γ -distances from the identity:

- **Hofer distance.** We set

$$d^H(\text{id}, \varphi; S_r) := \lim_{\varepsilon \rightarrow 0} d^H(\text{id}, \varphi_{H_{r,\varepsilon}}).$$

The above limit is well defined since if $H_{r,\varepsilon}$ and $K_{r,\varepsilon}$ are two different extensions of H in $\mathcal{H}_{r,\varepsilon}(H)$, then $\|H_{r,\varepsilon} - K_{r,\varepsilon}\|_{C^0} \leq 2\varepsilon$.

Note that $d^H(\text{id}, \varphi; S_r)$ only depends on S_r and φ , and not on H . We could also take a non-autonomous Hamiltonian, possibly non-convex, generating φ , provided it coincides with H near S_r . Indeed, the time one flow of $\varphi_{H_{r,\varepsilon}}^t = \varphi_H^{t f'_{r,\varepsilon}(H)}$ is determined by the knowledge of φ inside S_r and of φ_H^t in $S_{r+\varepsilon} \setminus S_r$. But the latter is determined by the hypersurfaces ∂S_s for $r \leq s \leq r + \varepsilon$. In any case we have the following lower bound: if $\varphi \neq \text{id}$ there is a ball in S_r such that $\varphi(B) \cap B = \emptyset$ and then $d^H(\text{id}, \varphi; S_r) \geq c(B) > 0$, where $c(B)$ is the Ekeland-Hofer capacity of B (cf. [8]).

- γ -**distance**. It can be defined as:

$$\gamma(\text{id}, \varphi; S_r) := \lim_{\varepsilon \rightarrow 0} \gamma(\text{id}, \varphi_{H_{r,\varepsilon}}).$$

Also this limit is well defined, for the same reasons as above. In fact, if $H_{r,\varepsilon}$ and $K_{r,\varepsilon}$ are two different extensions of H in $\mathcal{H}_{r,\varepsilon}(H)$, then

$$\gamma(\varphi_{H_{r,\varepsilon}}^1, \varphi_{K_{r,\varepsilon}}^1) \leq \|H_{r,\varepsilon} - K_{r,\varepsilon}\|_{C^0} \leq 2\varepsilon.$$

Similarly, $c_{\pm}(\varphi; S_r) := \lim_{\varepsilon \rightarrow 0} c_{\pm}(\varphi_{H_{r,\varepsilon}})$.

The same argument as before shows that $\gamma(\text{id}, \varphi; S_r)$ depends only on S_r and φ , and we have the same lower bound as for d^H . Analogously one can also define the associated asymptotic quantities d_{∞}^H , γ_{∞} and $c_{\pm,\infty}$.

We can now state our first result.

Theorem 1. *Let φ be the flow of an autonomous Tonelli Hamiltonian $H : \mathbb{T}^*\mathbb{T}^n \rightarrow \mathbb{R}$ and let $S_r \equiv \{(q, p) \in \mathbb{T}^*\mathbb{T}^n : H(q, p) \leq r\}$. Then, for each $r > \inf_{p \in \mathbb{R}^n} \alpha_H(p)$:*

$$\gamma_{\infty}(\text{id}, \varphi; S_r) = r + \beta_H(0).$$

More precisely, $c_{-,\infty}(\varphi; S_r) = \inf_{p \in \mathbb{R}^n} \overline{H}(p) = -\beta_H(0)$ and $c_{+,\infty}(\varphi; S_r) = r$.

Observe now that using Proposition 1, we obtain the following result.

Corollary 1. *Let φ be a Hamiltonian diffeomorphism generated by a Tonelli Hamiltonian $H : \mathbb{T}^*\mathbb{T}^n \rightarrow \mathbb{R}$. Then, for each $r > \inf_{p \in \mathbb{R}^n} \alpha_H(p)$:*

$$d_{\infty}^H(\text{id}, \varphi; S_r) \geq r + \beta_H(0).$$

The method used to prove Theorem 1 allows us to provide a new proof of Siburg's result in [15] (see also [13] for a generalization to general cotangent bundles). Let $B^*\mathbb{T}^n$ denote the unit ball cotangent bundle of \mathbb{T}^n , *i.e.* $B^*\mathbb{T}^n := \{(x, p) : \|p\| \leq 1\}$. Siburg considered the set of admissible Hamiltonians \mathcal{H}_S consisting of all smooth convex Hamiltonians $H : B^*\mathbb{T}^n \times \mathbb{T} \rightarrow \mathbb{R}$ that satisfy the following two conditions:

- S1) H vanishes on the boundary of $B^*\mathbb{T}^n$, *i.e.* $H(q, p, t) = 0$ if $\|p\| = 1$;
- S2) H admits a smooth extension $K_H : \mathbb{T}^*\mathbb{T}^n \times \mathbb{T} \rightarrow \mathbb{R}$ that is of Tonelli type and depends only on t and $\|p\|^2$ outside $B^*\mathbb{T}^n \times \mathbb{T}$.

Then, he defined the following group:

$$\mathcal{H}am_S(B^*\mathbb{T}^n) := \{\varphi : B^*\mathbb{T}^n \longrightarrow B^*\mathbb{T}^n \mid \varphi = \varphi_H^1 \text{ for some } H \in \mathcal{H}_S\}.$$

Observe that we cannot apply directly Theorem 1 in this setting, since these Hamiltonian diffeomorphisms are not necessarily generated by autonomous Hamiltonians. However, due to the special form of the Hamiltonians, the above results remain true.

Proposition 2. *Let $\varphi \in \mathcal{H}am_S(B^*\mathbb{T}^n)$. Then:*

$$d_\infty^H(\text{id}, \varphi) \geq \gamma_\infty(\text{id}, \varphi) = \beta_H(0).$$

More precisely, $c_{-, \infty}(\varphi) = \inf_{p \in \mathbb{R}} \overline{H}(p) = -\beta_H(0)$ and $c_{+, \infty}(\varphi) = 0$.

Remark 2. (i) The fact that $\beta_H(0)$ is independent of the extension K_H , has been proven in [15, Lemma 4.1]. Therefore, we define $\beta_H(0) := \beta_{K_H}(0)$. Observe that the independence could be also deduced from the above proposition and the fact that the γ -distance depends only on φ .

(ii) The inequality

$$d_\infty^H(\text{id}, \varphi) \geq \beta_H(0) \tag{1}$$

is due to Siburg [15, Theorem 5.1].

This lower bound (1) induced Siburg to ask the following question:

Question [15, page 94]: does equality hold in (1)?

We shall construct examples of convex Hamiltonian diffeomorphisms for which the asymptotic Hofer distance from the identity is strictly greater than the asymptotic γ -distance.

Theorem 2. *There exists $\varphi \in \mathcal{H}am_S(B^*\mathbb{T}^n)$, such that*

$$\gamma_\infty(\text{id}, \varphi) < d_\infty^H(\text{id}, \varphi).$$

An easy consequence of Theorem 2 and Proposition 2 is that the above question has a negative answer.

Corollary 2. *There exists $\varphi \in \mathcal{H}am_S(B^*\mathbb{T}^n)$ generated by convex Hamiltonian H , such that*

$$d_\infty^H(\text{id}, \varphi) > \beta_H(0).$$

4.1 Proof of Theorem 1

Proof of Theorem 1. Using [19, Theorem 4.2 (6)], we obtain

$$\begin{aligned} c_{-, \infty}(\varphi_H; S_r) &= \lim_{\varepsilon \rightarrow 0} c_{-, \infty}(\varphi_{H_r, \varepsilon}^1) = \\ &= \lim_{\varepsilon \rightarrow 0} \inf_{p \in \mathbb{R}^n} \overline{H}_{r, \varepsilon}(p). \end{aligned}$$

Similarly,

$$\begin{aligned} c_{+, \infty}(\varphi_H; S_r) &= \lim_{\varepsilon \rightarrow 0} c_{+, \infty}(\varphi_{H_r, \varepsilon}^1) = \\ &= \lim_{\varepsilon \rightarrow 0} \sup_{p \in \mathbb{R}^n} \overline{H}_{r, \varepsilon}(p). \end{aligned}$$

Observe that since $\|H_{r, \varepsilon} - H_r\|_{C^0} \leq \varepsilon$, then $\lim_{\varepsilon \rightarrow 0} \overline{H}_{r, \varepsilon} = \overline{H}_r$ uniformly. We want to prove that $\inf_{p \in \mathbb{R}^n} \overline{H}_r(p) = \inf_{p \in \mathbb{R}^n} \overline{H}(p)$. Clearly we have $\inf_p \overline{H}_r(p) \leq \inf_p \overline{H}(p)$, since $H_r \leq H$. We thus have to prove that $\inf_p \overline{H}_r(p) \geq \inf_p \overline{H}(p)$.

We shall need the following lemmata.

Lemma 1. *Let $f_{r, \varepsilon}$ be a function such that $f_{r, \varepsilon}(x) = x$ for $|x| \leq r$. Let us consider a Hamiltonian $H : \mathbb{T}^* \mathbb{T}^n \rightarrow \mathbb{R}$ and set $H_{r, \varepsilon} = f_{r, \varepsilon}(H)$. If there exists a Lagrangian Λ_p of cohomology class p , such that $\max_{\Lambda_p} H \leq r$, then $\overline{H}_{r, \varepsilon}(p) = \overline{H}(p)$.*

Proof. Since clearly the Poisson bracket of $H_{r, \varepsilon}$ and H vanishes, we have

$$\overline{H_{r, \varepsilon} - H} = \overline{H_{r, \varepsilon}} - \overline{H}$$

(see [19, Theorem 4.2 (7)]). It is thus enough to prove that $\overline{H_{r, \varepsilon}} - \overline{H}(p) = 0$. In other words, let K be a Hamiltonian vanishing on Λ_p , we must show that $\overline{K}(p) = 0$. But this follows immediately from Remark 4.3 in [19]. \square

Lemma 2. *Let f be a function in $C^0(\mathbb{R}, \mathbb{R})$ and $H : \mathbb{T}^* \mathbb{T}^n \rightarrow \mathbb{R}$ a convex superlinear Hamiltonian. Then, $\overline{f(H)} = f(\overline{H})$.*

Proof. Let us first consider the case in which f is smooth, convex and strictly increasing. Then, $f(H)$ is smooth, convex and superlinear and, using the characterization of the α -function given in [6], we obtain:

$$\overline{f(H)}(p) = \inf_{u \in C^\infty(\mathbb{T}^n)} \sup_{q \in \mathbb{T}^n} f(H(q, p + du(q))).$$

But clearly this coincides with $f(\inf_{u \in C^\infty(\mathbb{T}^n)} \sup_{q \in \mathbb{T}^n} H(q, p + du(q)) = f(\overline{H}(p))$. Now because $f(H)$ and $g(H)$ commute for any f and g , we have $\overline{f(H) - g(H)} = \overline{f(H)} - \overline{g(H)}$. Since differences of convex increasing functions are functions of bounded variations, and these form a dense subset among continuous functions², this concludes our proof. \square

As a result, we have that $\overline{H_r} = \overline{\min\{H, r\}} = \min\{\overline{H}, r\}$, hence $\inf_p \overline{H_r} = \inf_p \{\overline{H}, r\} = \inf_p \overline{H}(p)$, provided $r > \inf_p \overline{H}(p)$. And by the same argument, $\sup_p \overline{H_r} = \sup_p \min\{\overline{H}, r\} = r$ since $\lim_{\|p\| \rightarrow \infty} \overline{H}(p) = +\infty$. This concludes the proof of Theorem 1. \square

4.2 Proof of Proposition 2

Before entering into the details of the proof, let us observe that also in this case $\varphi \in \mathcal{H}am_S(B^*\mathbb{T}^n)$ is not necessarily compactly supported, therefore we need to define compactly supported extensions. Let us denote $B_r^*\mathbb{T}^n := \{(q, p) : \|p\| \leq r\}$. If H is a generating Hamiltonian for φ , for each ε sufficiently small we can consider new compactly supported Hamiltonians $H_\varepsilon = f_{0,\varepsilon}(K_H)$ supported in a neighborhood of $B^*\mathbb{T}^n \times \mathbb{T}$. The Hamiltonian H_ε has the following properties (see section 4 for the properties of $f_{0,\varepsilon}$):

- (1) it coincides with H in $B^*\mathbb{T}^n \times \mathbb{T}$;
- (2) it is non-negative and bounded by ε outside $B^*\mathbb{T}^n \times \mathbb{T}$ and satisfy $\sqrt{\varepsilon} \|d_p H_\varepsilon\| \leq C$ outside $B^*\mathbb{T}^n \times \mathbb{T}$, for some constant C depending only on H ;
- (3) it depends only on $\|p\|^2$ and t outside $B^*\mathbb{T}^n \times \mathbb{T}$.

We denote the set of all these Hamiltonians by $\mathcal{H}_\varepsilon(H)$. As before, we set

$$\begin{aligned} d^H(\text{id}, \varphi) &:= \lim_{\varepsilon \rightarrow 0} d^H(\text{id}, \varphi_{H_\varepsilon}) \\ \gamma(\text{id}, \varphi) &:= \lim_{\varepsilon \rightarrow 0} \gamma(\text{id}, \varphi_{H_\varepsilon}) \\ c_\pm(\varphi) &:= \lim_{\varepsilon \rightarrow 0} c_\pm(\varphi_{H_\varepsilon}). \end{aligned}$$

For the same reasons as before, these definitions are independent of the chosen H_ε and one can define analogously the associated asymptotic quantities.

²This is true by definition without the convexity requirement. However if $h = f - g$ we may also write $h(x) = (f(x) + cx^2) - (g(x) + cx^2)$, and for c sufficiently large both terms are convex.

Proof of Proposition 2. Note first that the left-hand side inequality is just Proposition 1. Let $\varphi \in \mathcal{H}am_S(B^*\mathbb{T}^n)$ be generated by a convex Hamiltonian H and let K_H denote its smooth extension to $\mathbb{T}^*\mathbb{T}^n \times \mathbb{T}$, given by condition (S2) in the definition of \mathcal{H}_S . Then,

$$\begin{aligned} \beta_H(0) &:= \beta_{K_H}(0) = \sup_{c \in H^1(\mathbb{T}^n; \mathbb{R})} (\langle c, 0 \rangle - \alpha_{K_H}(c)) = - \inf_{c \in H^1(\mathbb{T}^n; \mathbb{R})} \alpha_{K_H}(c) = \\ &= - \inf_{p \in \mathbb{R}^n} \overline{K}_H(p), \end{aligned} \quad (2)$$

where in the last equality we used that α_{K_H} coincides with the symplectic homogenization \overline{K}_H .

Let \mathcal{B}_1 denote the closed unit ball in \mathbb{R}^n with the standard norm. First, we want to prove that $\overline{K}_H \equiv 0$ on $\partial\mathcal{B}_1$. Then, using the convexity of \mathcal{B}_1 , it follows easily that $\overline{K}_H \leq 0$ in \mathcal{B}_1 and that:

$$\max_{\mathcal{B}_1} \overline{K}_H = 0 \quad \text{and} \quad \min_{\mathcal{B}_1} \overline{K}_H = \min_{p \in \mathbb{R}^n} \overline{K}_H.$$

In order to prove that \overline{K}_H vanishes on $\partial\mathcal{B}_1$, observe that for each $p_0 \in \partial\mathcal{B}_1$, K_H vanishes on the Lagrangian submanifold $\Lambda_{p_0} := \mathbb{T}^n \times \{p_0\}$. The claim then follows from [19, Theorem 4.2 (5)].

The proof of Proposition 2 can now be obtained applying [19, Theorem 4.2 (6)].

In fact,

$$\begin{aligned} \gamma_\infty(\text{id}, \varphi) &= \lim_{\varepsilon \rightarrow 0} \gamma_\infty(\text{id}, \varphi_{H_\varepsilon}) = \\ &= \lim_{\varepsilon \rightarrow 0} c_{+, \infty}(\varphi_{H_\varepsilon}) - \lim_{\varepsilon \rightarrow 0} c_{-, \infty}(\varphi_{H_\varepsilon}) = \\ &= \lim_{\varepsilon \rightarrow 0} \sup_{p \in \mathbb{R}^n} \overline{H}_\varepsilon(p) - \lim_{\varepsilon \rightarrow 0} \inf_{p \in \mathbb{R}^n} \overline{H}_\varepsilon(p). \end{aligned} \quad (3)$$

Using the fact that K_H and H_ε coincide on $B^*\mathbb{T}^n$ and proceeding as in Lemma 1, one can deduce that $\overline{H}_\varepsilon(p) = \overline{K}_H(p)$ for each $p \in \mathcal{B}_1$. Moreover, since H_ε is bounded by ε , we obtain:

$$0 = \sup_{p \in \mathcal{B}_1} \overline{K}_H(p) = \sup_{p \in \mathcal{B}_1} \overline{H}_\varepsilon(p) \leq \lim_{\varepsilon \rightarrow 0} \sup_{p \in \mathbb{R}^n} \overline{H}_\varepsilon(p) \leq \lim_{\varepsilon \rightarrow 0} \varepsilon = 0.$$

Similarly for the infimum of \overline{H}_ε , we have

$$0 \geq \inf_{p \in \mathcal{B}_1} \overline{K}_H(p) = \inf_{p \in \mathcal{B}_1} \overline{H}_\varepsilon(p) = \lim_{\varepsilon \rightarrow 0} \inf_{p \in \mathcal{B}_1} \overline{H}_\varepsilon(p),$$

using for the last equality the fact that the left-hand side is independent of ε . Now we use the fact that $H_\varepsilon \geq 0$ on $\mathbb{T}^n \times \{p\}$ for any $p \in \mathbb{R}^n \setminus \mathcal{B}_1$ and [19, Theorem 4.2 (5)], to conclude that $\overline{H}_\varepsilon(p) \geq 0$ for $p \in \mathbb{R}^n \setminus \mathcal{B}_1$, and hence $\inf_{p \in \mathbb{R}^n} \overline{H}_\varepsilon(p) = \inf_{p \in \mathcal{B}_1} \overline{H}_\varepsilon(p)$.

Therefore, substituting in (3) and using (2) we can conclude:

$$\begin{aligned} \gamma_\infty(\text{id}, \varphi) &= \limsup_{\varepsilon \rightarrow 0} \inf_{\mathbb{R}^n} \overline{H}_\varepsilon(p) - \liminf_{\varepsilon \rightarrow 0} \inf_{\mathbb{R}^n} \overline{H}_\varepsilon(p) = \\ &= - \inf_{p \in \mathcal{B}_1} \overline{K}_H(p) = \beta_H(0). \end{aligned}$$

□

4.3 Proof of Theorem 2

In this section we shall construct examples of $\varphi \in \mathcal{H}am_S(B^*\mathbb{T}^n)$ generated by a convex Hamiltonian H , such that $\gamma_\infty(\text{id}, \varphi) < d_\infty^H(\text{id}, \varphi)$.

Proof of Theorem 2. The basic observation is that in the compactly supported case the Hofer distance can be bounded from below in terms of the so-called *Calabi invariant* (see [4] and [3]):

$$\mathcal{C}al(\varphi) := \int_0^1 \int_{\mathbb{T}^*\mathbb{T}^n} H(q, p, t) \omega^n dt.$$

This invariant only depends on φ and not on the path defined by H . Indeed, let us consider the Liouville form $\lambda = \sum_{j=1}^n p_j dq_j$ on $\mathbb{T}^*\mathbb{T}^n$. Then $\varphi^*\lambda - \lambda = df_\varphi$, and in the compactly supported case, an easy computation (see [3] for instance) shows that

$$\mathcal{C}al(\varphi) = \frac{1}{n+1} \int_{\mathbb{T}^*\mathbb{T}^n} f_\varphi \omega^n. \quad (4)$$

Let us adapt this to the case of $\mathcal{H}am_S(B^*\mathbb{T}^n)$. In our situation we set again

$$\mathcal{C}al(\varphi) = \int_0^1 \int_{B^*\mathbb{T}^n} H(q, p, t) \omega^n dt.$$

Note that if H_ε is a compactly supported extension of H (as defined in section 4.2), we have

$$\lim_{\varepsilon \rightarrow 0} \mathcal{C}al(\varphi_{H_\varepsilon}) = \mathcal{C}al(\varphi). \quad (5)$$

Indeed, let us denote by $B_{\rho(\varepsilon)}^* \mathbb{T}^n$ a disc bundle of radius $\rho(\varepsilon)$ that contains the support of $H_\varepsilon - \varepsilon$ (i.e. contains $K_H^{-1}(\varepsilon)$), and notice that since $d_p K_H(q, p, t) \cdot p$ is non-zero near $\partial B^* \mathbb{T}^n$ we have that $\rho(\varepsilon) - 1 = O(\varepsilon)$. Applying formula (4) to φ_{H_ε} instead of φ , we get

$$\mathcal{Cal}(\varphi_{H_\varepsilon}) = \frac{1}{n+1} \int_{B^* \mathbb{T}^n} f_\varphi + \frac{1}{n+1} \int_{B_{\rho(\varepsilon)}^* \mathbb{T}^n \setminus B^* \mathbb{T}^n} f_{\varphi_{H_\varepsilon}}; \quad (6)$$

but on $B_{\rho(\varepsilon)}^* \mathbb{T}^n \setminus B^* \mathbb{T}^n$, $f_{\varphi_{H_\varepsilon}}$ is given by $\langle d_p H_\varepsilon(q, p, t), p \rangle - H_\varepsilon(q, p, t)$. It is thus enough to check that the integral of this quantity on $(B_{\rho(\varepsilon)}^* \mathbb{T}^n \setminus B^* \mathbb{T}^n) \times \mathbb{T}$ goes to zero as ε goes to zero, or else that

$$\varepsilon \|\langle dH_\varepsilon(q, p, t), p \rangle - H_\varepsilon(q, p, t)\|_{C^0} \xrightarrow{\varepsilon \rightarrow 0} 0.$$

Since $\|H_\varepsilon\|_{C^0}$ and $\varepsilon \|d_p H_\varepsilon(q, p, t)\|_{C^0}$ go to zero on $(B_{\rho(\varepsilon)}^* \mathbb{T}^n \setminus B^* \mathbb{T}^n) \times \mathbb{T}$, this clearly holds.

We now compare the Hofer distance with the Calabi invariant.

Lemma 3. *Let $\varphi \in \mathcal{Ham}_S(B^* \mathbb{T}^n)$. Then,*

$$d_\infty^H(\text{id}, \varphi) \geq \frac{1}{\text{Vol}(B^* \mathbb{T}^n)} |\mathcal{Cal}(\varphi)|.$$

Proof. Let $\varphi \in \mathcal{Ham}_S(B^* \mathbb{T}^n)$ and $H(q, p, t)$ be a (not necessarily convex) generating Hamiltonian. Recall from the definition of $\mathcal{Ham}_S(B^* \mathbb{T}^n)$ that $H(q, p, t) = 0$ on $\|p\| = 1$ and that it admits a smooth extension $H_\varepsilon : \mathbb{T}^* \mathbb{T}^n \times \mathbb{T} \rightarrow \mathbb{R}$ in $\mathcal{H}_\varepsilon(H)$, which is a function only of t and $\|p\|^2$ outside $B^* \mathbb{T}^n \times \mathbb{T}$ and which is bounded by ε outside $B^* \mathbb{T}^n \times \mathbb{T}$ (see section 4.2).

Now, denoting by φ_ε the time-one flow $\varphi_{H_\varepsilon}^1$, we may write:

$$\begin{aligned} d^H(\text{id}, \varphi_\varepsilon) &= \int_0^1 \text{Osc}_{B_{\rho(\varepsilon)}^* \mathbb{T}^n}(H_{\varepsilon, t}) dt = \\ &= \frac{1}{\text{Vol}(B_{\rho(\varepsilon)}^* \mathbb{T}^n)} \int_0^1 \int_{\mathbb{T}^* \mathbb{T}^n} \text{Osc}_{\mathbb{T}^* \mathbb{T}^n}(H_{\varepsilon, t}) \omega^n dt \geq \\ &\geq \frac{1}{\text{Vol}(B_{\rho(\varepsilon)}^* \mathbb{T}^n)} \left| \int_0^1 \int_{\mathbb{T}^* \mathbb{T}^n} H_\varepsilon(q, p, t) \omega^n dt \right| = \\ &= \frac{1}{\text{Vol}(B_{\rho(\varepsilon)}^* \mathbb{T}^n)} |\mathcal{Cal}(\varphi_\varepsilon)|. \end{aligned}$$

Hence, using (5), we obtain:

$$\begin{aligned}
d^H(\text{id}, \varphi) &= \lim_{\varepsilon \rightarrow 0} d^H(\text{id}, \varphi_\varepsilon) \geq \\
&\geq \lim_{\varepsilon \rightarrow 0} \frac{1}{\text{Vol}(B_{\rho(\varepsilon)}^* \mathbb{T}^n)} |\mathcal{Cal}(\varphi_\varepsilon)| = \\
&= \frac{1}{\text{Vol}(B^* \mathbb{T}^n)} |\mathcal{Cal}(\varphi)|.
\end{aligned}$$

Since $H(q, p, t+1) = H(q, p, t)$, we can conclude that:

$$\begin{aligned}
d_\infty^H(\text{id}, \varphi) &\geq \frac{1}{\text{Vol}(B^* \mathbb{T}^n)} \lim_{k \rightarrow +\infty} \left(\frac{1}{k} |\mathcal{Cal}(\varphi_H^k)| \right) = \\
&= \frac{1}{\text{Vol}(B^* \mathbb{T}^n)} \left| \int_0^1 \int_{B^* \mathbb{T}^n} H(t, q, p) \omega^n dt \right| = \\
&= \frac{1}{\text{Vol}(B^* \mathbb{T}^n)} |\mathcal{Cal}(\varphi)|.
\end{aligned}$$

□

In order to find our example of $\varphi \in \mathcal{H}am_S(B^* \mathbb{T}^n)$ such that $\gamma_\infty(\text{id}, \varphi) < d_\infty^H(\text{id}, \varphi)$, it is sufficient to find φ such that

$$\gamma_\infty(\text{id}, \varphi) < \frac{1}{\text{Vol}(B^* \mathbb{T}^n)} |\mathcal{Cal}(\varphi)|.$$

Let U_δ be a cube of side $\delta < \frac{1}{3}$ contained in \mathbb{T}^n . Let H be a negative convex Hamiltonian of the form $H(q, p) = \gamma(q)(\|p\|^2 - 1)$, such that $\gamma(q) \geq 0$ and

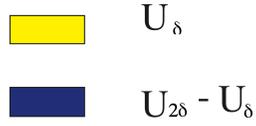
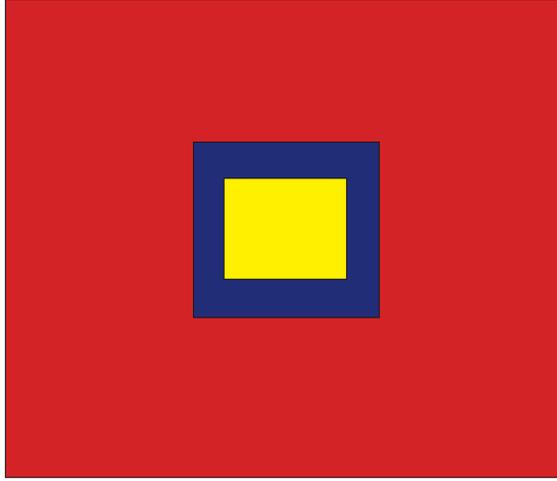


Figure 1: The sets $U_\delta, U_{2\delta}$

$$\gamma(q) = \begin{cases} C & \text{on } U_\delta \\ c & \text{on } \mathbb{T}^n \setminus U_{2\delta}. \end{cases}$$

where $c \ll C$.

Observe that:

$$|\mathcal{Cal}(\varphi_H^1)| := \lim_{\varepsilon \rightarrow 0} |\mathcal{Cal}(\varphi_{H_\varepsilon}^1)| \geq [\delta^n C + c(1 - 2^n \delta^n)] \int_{\{\|p\| \leq 1\}} (1 - \|p\|^2) dp.$$

If we set $k := \int_{\{\|p\| \leq 1\}} (1 - \|p\|^2) dp$, then

$$|\mathcal{Cal}(\varphi_H^1)| \geq [\delta^n C + c(1 - 2^n \delta^n)]k.$$

In order to conclude the proof, it is sufficient to prove that

$$\lim_{\varepsilon \rightarrow 0} |\overline{H}_\varepsilon(p)| < \frac{\delta^n Ck}{\text{Vol}(B^*\mathbb{T}^n)} \quad \text{for all } p,$$

where as usual \overline{H}_ε denotes the symplectic homogenization of H_ε . Then, applying Proposition 2 and Lemma 3, it follows that

$$\gamma_\infty(\text{id}, \varphi) < \frac{1}{\text{Vol}(B^*\mathbb{T}^n)} |\text{Cal}(\varphi)| \leq d_\infty^H(\varphi, \text{id}).$$

In order to prove this, we use the fact that if there exists a Lagrangian submanifold $\Lambda_p = \psi(\mathbb{T}^n \times \{p\})$, where $\psi \in \mathcal{H}am(\mathbb{T}^*\mathbb{T}^n)$, such that $-H_\varepsilon|_{\Lambda_p} < A$, then $-\overline{H}_\varepsilon(p) \leq A$; see [19, Theorem 3.2]. We look for Λ_p in the form

$$\Lambda_p = \{(q, p + df(q)) : q \in \mathbb{T}^n\},$$

with f satisfying the condition $\|p + df(q)\| = 1$ on $U_{2\delta}$.

Lemma 4. *For all vectors p in \mathbb{R}^n there exists a smooth function f on \mathbb{T}^n such that $\|df(q) + p\| = 1$ on $U_{2\delta}$.*

Proof. We must find a vector field $u(q)$ of norm 1 on $U_{2\delta}$, such that $df(q) = u(q) - p$ on $U_{2\delta}$. Take u to be constant on $U_{2\delta}$, then $f(q) = \langle u - p, q \rangle$ and extend this to a smooth function on \mathbb{T}^n . \square

Then,

$$\begin{aligned} -H_\varepsilon(q, p + df(q)) &= 0 \quad \text{on } U_{2\delta} \quad (\text{because } \|df(q) + p\| = 1) \\ -\varepsilon \leq -H_\varepsilon(q, p + df(q)) &\leq c \quad \text{on } \mathbb{T}^n \setminus U_{2\delta}. \end{aligned}$$

Hence $-\overline{H}(p) = |\overline{H}(p)| = \lim_{\varepsilon \rightarrow 0} |\overline{H}_\varepsilon(p)| \leq c$ for all p . Therefore, we proved that

$$\gamma_\infty(\varphi, \text{id}) = \lim_{\varepsilon \rightarrow 0} \text{Osc}(\overline{H}_\varepsilon) \leq \lim_{\varepsilon \rightarrow 0} \sup_{p \in \mathbb{R}^n} |\overline{H}_\varepsilon(p)| \leq c$$

Provided we choose our constants to satisfy $c < \frac{\delta^n Ck}{\text{Vol}(B^*\mathbb{T}^n)}$, this concludes the proof of Theorem 2. \square

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